

Random quantum channels - graphical calculus -

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joint work with Benot Collins

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Random quantum channels
&
additivity problems

Additivity for MOE of quantum channels

- **Quantum channels:** CPTP maps $\Phi : \mathcal{M}_{\text{in}}(\mathbb{C}) \rightarrow \mathcal{M}_{\text{out}}(\mathbb{C})$.

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- **NO !!!**

- $p > 1$: Hayden '07, Hayden + Winter '08
- $p = 1$: Hastings '09

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$$\Phi(\rho) = \text{Tr}_{\text{aux}}(V\rho V^*),$$

where V is a Haar partial isometry

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- Equivalently, via the Stinespring dilation theorem

$$\Phi(\rho) = \text{Tr}_{\text{aux}}(U(\rho \otimes P_y)U^*),$$

where $y \in \mathbb{C}^{\frac{\text{out} \times \text{aux}}{\text{in}}}$ and $U \in \mathcal{M}_{\text{out} \times \text{aux}}(\mathbb{C})$ is a Haar unitary matrix.

Choice of parameters

- $\text{in} = tnk$,
- $\text{out} = k$,
- $\text{aux} = n$,

where $n, k \in \mathbb{N}$ and $t \in (0, 1)$. In general, we shall assume that

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- k is fixed, but “large”;
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We are thus considering random channels

$$\begin{aligned}\Phi : \mathcal{M}_{tnk}(\mathbb{C}) &\rightarrow \mathcal{M}_k(\mathbb{C}) \\ \rho &\mapsto \text{Tr}_n[U(\rho \otimes P_y)U^*],\end{aligned}$$

where $y \in \mathbb{C}^{t-1}$ is fixed (and irrelevant) and $U \in \mathcal{U}(nk)$ is a Haar random unitary matrix.

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Strategy

- Use trivial bound

$$H_{\min}^p(\Phi \otimes \overline{\Phi}) \leq H^p([\Phi \otimes \overline{\Phi}](X_{12})),$$

for a particular choice of $X_{12} \in \mathcal{M}_{tnk}(\mathbb{C}) \otimes \mathcal{M}_{tnk}(\mathbb{C})$.

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- Bound entropies of the (random) density matrix

$$Z = [\Phi \otimes \bar{\Phi}](E_{tnk}) \in \mathcal{M}_{k^2}(\mathbb{C}).$$

Main result for product channel

Theorem (Collins + N. '09)

For all k, t , almost surely as $n \rightarrow \infty$, the eigenvalues of $Z = [\Phi \otimes \bar{\Phi}](E_{tnk})$ converge to

$$\left(t + \frac{1-t}{k^2}, \underbrace{\frac{1-t}{k^2}, \dots, \frac{1-t}{k^2}}_{k^2-1 \text{ times}} \right).$$

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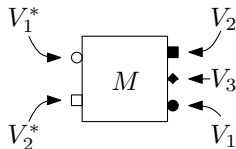
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- Two improvements:
 - ① “better” largest eigenvalue,
 - ② knowledge of the whole spectrum.
- However, smaller eigenvalues are the “worst possible”.
- Precise knowledge of eigenvalue \rightsquigarrow **optimal** estimates for entropies.

Graphical calculus for random quantum channels

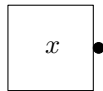
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Boxes & wires

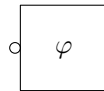
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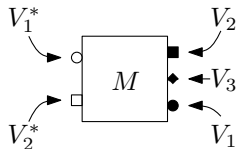
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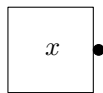
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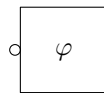
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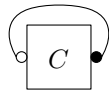
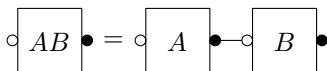


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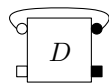


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- Tensor contractions (or traces) $V \otimes V^* \rightarrow \mathbb{C} \rightsquigarrow$ wires.



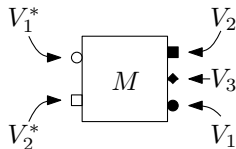
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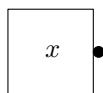
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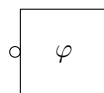
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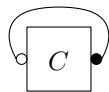
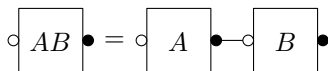


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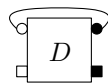


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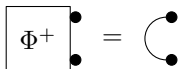


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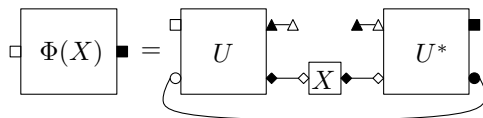
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- Bell state $\Phi^+ = \sum_{i=1}^{\dim V_1} e_i \otimes e_i \in V_1 \otimes V_1$

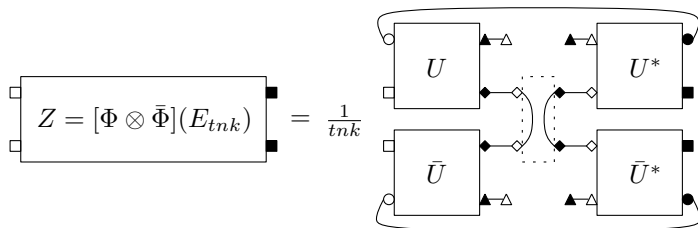


Graphical representation of quantum channels

- Single channel



- Product of conjugate channels



- Decorations/labels

$$\circ = \mathbf{C}^n$$

$$\square = \mathbf{C}^k$$

$$\diamond = \mathbf{C}^{tnk}$$

$$\blacktriangle = \mathbf{C}^{t^{-1}}$$

Proof strategy for a.s. spectrum of random channels

- Use the [method of moments](#)

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① Convergence in moments:

$$\mathbb{E} \operatorname{Tr}(Z^p) \rightarrow \left(t + \frac{1-t}{k^2} \right)^p + (k^2 - 1) \left(\frac{1-t}{k^2} \right)^p ;$$

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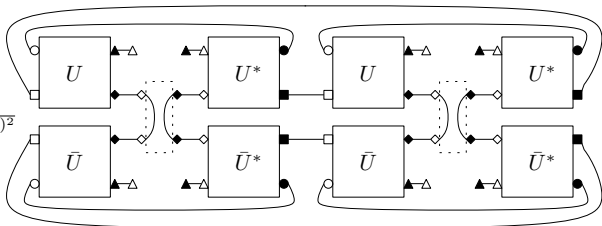
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- Example**

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- Using matrix coordinates, we can reduce our problem to computing integrals over the unitary group.

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Theorem (Weingarten formula)

Let d be a positive integer and $\mathbf{i} = (i_1, \dots, i_p)$, $\mathbf{i}' = (i'_1, \dots, i'_p)$, $\mathbf{j} = (j_1, \dots, j_p)$, $\mathbf{j}' = (j'_1, \dots, j'_p)$ be p -tuples of positive integers from $\{1, 2, \dots, d\}$. Then

$$\int_{\mathcal{U}(d)} U_{i_1 j_1} \cdots U_{i_p j_p} \overline{U_{i'_1 j'_1}} \cdots \overline{U_{i'_p j'_p}} dU = \sum_{\alpha, \beta \in \mathcal{S}_p} \delta_{i_1 i'_{\alpha(1)}} \cdots \delta_{i_p i'_{\alpha(p)}} \delta_{j_1 j'_{\beta(1)}} \cdots \delta_{j_p j'_{\beta(p)}} \text{Wg}(d, \alpha \beta^{-1}).$$

If $p \neq p'$ then

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- There is a [graphical](#) way of reading this formula on the diagrams !

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- 4 For all $i = 1, \dots, p$, add a wire between each white decoration of the i -th U box and the corresponding white decoration of the $\alpha(i)$ -th \overline{U} box. In a similar manner, use β to pair black decorations.

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- 5 Erase all U and \overline{U} boxes. The resulting diagram is denoted by $\mathcal{D}_{(\alpha, \beta)}$.

Theorem

$$\mathbb{E}\mathcal{D} = \sum_{\alpha, \beta} \mathcal{D}_{(\alpha, \beta)} \text{Wg}(d, \alpha\beta^{-1}).$$

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- $\mathcal{D}_{(\alpha, \beta)}$ is a collection of loops associated to vector spaces of dimensions n , k and tnk .

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- One needs to compute the contribution of each diagram $\mathcal{D}_{(\alpha, \beta)}$, where $\alpha, \beta \in \mathcal{S}_{2p}$.
- $\mathcal{D}_{(\alpha, \beta)}$ is a collection of loops associated to vector spaces of dimensions n , k and tnk .
- After doing the loop combinatorics, one is left with maximizing over S_{2p}^2 quantities such as

$$\#(\gamma^{-1}\alpha) + \#(\alpha^{-1}\beta) + \#(\beta^{-1}\delta),$$

where γ and δ are permutations coding the initial wiring of U/\bar{U} boxes and $\#(\cdot)$ is the number of cycles function.

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- Geodesic problems in symmetric groups \Rightarrow non-crossing partitions \Rightarrow free probability.
- Asymptotic for Weingarten weights:

$$\text{Wg}(d, \sigma) = d^{-(p+|\sigma|)}(\text{Mob}(\sigma) + O(d^{-2})).$$

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- Improved bounds for MOE of product channels
- Other applications to QIT (work in progress with B. Collins and K. Życzkowski)

Thank you !

Next talk \rightsquigarrow bounds for 1 channel

<http://arxiv.org/abs/0905.2313>

<http://arxiv.org/abs/0906.1877>