

# Some applications of free spectrahedra to quantum information theory

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# What this talk is about

We establish a connection between two topics:

- **Incompatibility** of quantum measurements
- Inclusion of **free spectrahedra**, geometric objects which generalize polyhedra

The relation we put forward allows us to find the **most incompatible** quantum measurements in different settings. These objects are of great importance for the **foundations of quantum mechanics** due to their non-classical behavior. They also encode, in some sense, the limits of quantum theory.

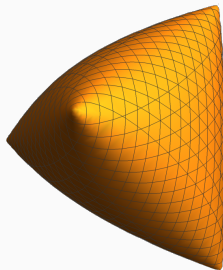
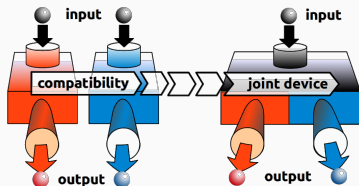
# Talk outline

Incompatibility in QM

Free spectrahedra

Main results

Proof ideas



# Incompatibility in QM

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# Quantum measurements. Compatibility

- Quantum states  $\rightsquigarrow$  **unit vectors in a Hilbert space**:  $\psi \in \mathbb{C}^d$ ,  $\|\psi\| = 1$
- Quantum measurements are modeled by **POVMs** (Positive Operator Valued Measures)

$$A_1, \dots, A_k \in M_d(\mathbb{C})^{\text{sa}}, \quad A_i \geq 0, \quad \sum_{i=1}^k A_i = I_d$$

- Born's rule**:  $\mathbb{P}[\text{outcome } i \text{ is observed}] = \langle \psi, A_i \psi \rangle$

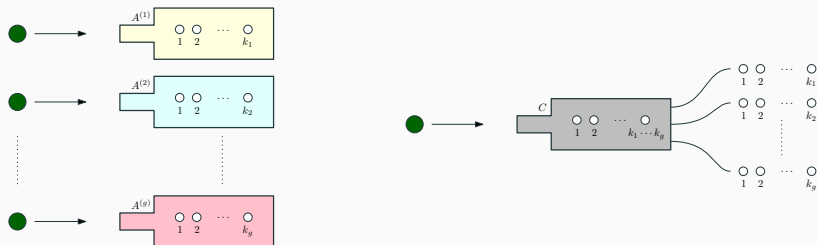
## Definition

Two POVMs,  $A = (A_1, \dots, A_k)$  and  $B = (B_1, \dots, B_l)$ , are called **compatible** if there exists a third POVM  $C = (C_{ij})_{i \in [k], j \in [l]}$  such that

$$\forall i \in [k], \quad A_i = \sum_{j=1}^l C_{ij} \quad \text{and} \quad \forall j \in [l], \quad B_j = \sum_{i=1}^k C_{ij}$$

The definition generalizes to  $g$ -tuples of POVMs  $A^{(1)}, \dots, A^{(g)}$ , having respectively  $k_1, \dots, k_g$  outcomes, where the **joint** POVM  $C$  has outcome set  $[k_1] \times \dots \times [k_g]$ .

# What does it mean?



- Compatible measurements can be simulated by a single joint measurement, by **classically post-processing** its outputs
- Examples:
  - ① **Trivial** POVMs  $A = (p_i I_d)$  and  $B = (q_j I_d)$  are compatible
  - ② **Commuting** POVMs  $[A_i, B_j] = 0$  are compatible
  - ③ If the POVM  $A$  is **projective**, then  $A$  and  $B$  are compatible iff they commute
- Compatibility appears also in the setting of quantum channels [HMZ16]

# Noisy POVMs

- POVMs can be made compatible by adding **noise**, i.e. mixing in trivial POVMs
- Example: dichotomic POVMs and white noise,  $s \in [0, 1]$

$$(E, I - E) \mapsto s(E, I - E) + (1 - s)\left(\frac{I}{2}, \frac{I}{2}\right) \quad \text{or} \quad E \mapsto sE + (1 - s)\frac{I}{2}$$

- Taking  $s = 1/2$  suffices to render any pair of dichotomic POVMs compatible [DFK19]

$$\rightsquigarrow \text{define } C_{ij} := (E_i + F_j)/4$$

- From now on, we focus on dichotomic (YES/NO) POVMs

## Definition

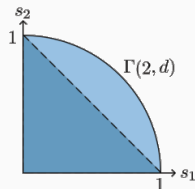
The **compatibility region** for  $g$  measurements on  $\mathbb{C}^d$  is the set

$$\Gamma(g, d) := \{s \in [0, 1]^g : \text{for all quantum effects } E_1, \dots, E_g \in \mathcal{M}_d(\mathbb{C}), \\ \text{the noisy versions } s_i E_i + (1 - s_i)I_d/2 \text{ are compatible}\}$$

# Compatibility region

$\Gamma(g, d) := \{s \in [0, 1]^g : \text{for all quantum effects } E_1, \dots, E_g \in \mathcal{M}_d(\mathbb{C}),$   
the noisy versions  $s_i E_i + (1 - s_i) I_d / 2$  are compatible}

- The set  $\Gamma(g, d)$  is convex
- For all  $i \in [g]$ ,  $e_i \in \Gamma(g, d)$ : every measurement is compatible with  $g - 1$  trivial measurements
- For  $d \geq 2$ ,  $(1, 1, \dots, 1) \notin \Gamma(g, d)$ : there exist incompatible measurements
- For all  $d \geq 2$ ,  $\Gamma(2, d)$  is a quarter-circle [BH08, BHSS13]
- Generally speaking, the set  $\Gamma(g, d)$  tells us how robust (to noise) is the incompatibility of  $g$  dichotomic measurements on  $\mathbb{C}^d$



**GOAL:** compute the set  $\Gamma(g, d)$  for all  $g, d$



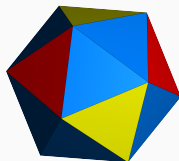
# Free spectrahedra

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# Free spectrahedra

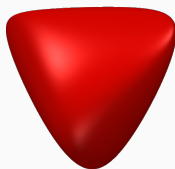
- A **polyhedron** is defined as the intersection of half-spaces

$$\{x \in \mathbb{R}^g : \langle h_i, x \rangle \leq 1, \quad \forall i \in [k]\}$$



- A **spectrahedron** is given by PSD constraints:  
for  $A = (A_1, \dots, A_g) \in (\mathcal{M}_d^{sa})^g$

$$\mathcal{D}_A(1) := \{x \in \mathbb{R}^g : \sum_{i=1}^g x_i A_i \leq I_d\}$$



- $\mathcal{D}_{(\sigma_X, \sigma_Y, \sigma_Z)}(1) = \{(x, y, z) \in \mathbb{R}^3 : x\sigma_X + y\sigma_Y + z\sigma_Z \leq I_2\} = \text{Bloch ball}$
- A **free spectrahedron** is the matricization of a spectrahedron [Vin14]

$$\mathcal{D}_A := \bigsqcup_{n=1}^{\infty} \mathcal{D}_A(n) \quad \text{with} \quad \mathcal{D}_A(n) := \{X \in (\mathcal{M}_n^{sa})^g : \sum_{i=1}^g X_i \otimes A_i \leq I_{nd}\}$$

## Examples: the cube and the diamond

The **matrix cube** is the free spectrahedron defined by

$$\mathcal{D}_{\square,g} := \bigcap_{n=1}^{\infty} \{X \in (\mathcal{M}_n^{sa})^g : \|X_i\| \leq 1, \quad \forall i \in [g]\}$$

- At level one,  $\mathcal{D}_{\square,g}(1)$  is the unit ball of the  $\ell^\infty$  norm on  $\mathbb{R}^g$
- As a free spectrahedron, it is defined by  $2g \times 2g$  diagonal matrices  $\mathcal{D}_{\square,g} = \mathcal{D}_{K_1, \dots, K_g}$ , with  $K_i = \text{diag}(e_i) \oplus \text{diag}(-e_i)$

The **matrix diamond** is the free spectrahedron defined by

$$\mathcal{D}_{\diamond,g} := \bigcap_{n=1}^{\infty} \{X \in (\mathcal{M}_n^{sa})^g : \sum_{i=1}^g \varepsilon_i X_i \leq I_n, \quad \forall \varepsilon \in \{\pm 1\}^g\}$$

- At level one,  $\mathcal{D}_{\diamond,g}(1)$  is the unit ball of the  $\ell^1$  norm on  $\mathbb{R}^g$
- As a free spectrahedron, it is defined by  $2^g \times 2^g$  diagonal matrices  $\mathcal{D}_{\diamond,g} = \mathcal{D}_{L_1, \dots, L_g}$ , with  $L_i = I_2 \otimes \dots \otimes I_2 \otimes \text{diag}(1, -1) \otimes I_2 \otimes \dots \otimes I_2$

# Spectrahedral inclusion

- Consider two free spectrahedra defined by  $g$ -tuples of matrices  $(A_1, \dots, A_g)$  and  $(B_1, \dots, B_g)$
- We say that  $\mathcal{D}_A$  is **contained** in  $\mathcal{D}_B$  and we write  $\mathcal{D}_A \subseteq \mathcal{D}_B$  if, for all  $n \geq 1$ ,  $\mathcal{D}_A(n) \subseteq \mathcal{D}_B(n)$
- Clearly,  $\mathcal{D}_A \subseteq \mathcal{D}_B \implies \mathcal{D}_A(1) \subseteq \mathcal{D}_B(1)$ . For the converse implication to hold, one may need to shrink  $\mathcal{D}_A$ ...

## Definition

For a free spectrahedron  $\mathcal{D}_A$ , we define its set of **inclusion constants** as

$$\Delta_A(g, d) := \{s \in [0, 1]^g : \text{for all } g\text{-tuples } B_1, \dots, B_g \in \mathcal{M}_d(\mathbb{C})^{\text{sa}}, \\ \mathcal{D}_A(1) \subseteq \mathcal{D}_B(1) \implies s \cdot \mathcal{D}_A \subseteq \mathcal{D}_B\}$$

- The inclusion constants for the matrix cube play an important role in combinatorial optimization [HKMS19]
- We shall be concerned with the inclusion set for the **matrix diamond**, which we denote by  $\Delta(g, d)$ ; it is the same as the one of the matrix cube [BN22a]

## **Main results**

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# Compatibility in QM $\iff$ matrix diamond inclusion

To a  $g$ -tuple of selfadjoint matrices  $E \in (\mathcal{M}_d^{sa})^g$ , we associate the free spectrahedron defined by the matrices  $2E_i - I_d$ :

$$\mathcal{D}_{2E-I} := \bigcup_{n=1}^{\infty} \{X \in (\mathcal{M}_n^{sa})^g : \sum_{i=1}^g X_i \otimes (2E_i - I_d) \leq I_{nd}\}$$

## Theorem ([BN18])

Let  $E \in (\mathcal{M}_d^{sa})^g$  be  $g$ -tuple of selfadjoint matrices. Then:

- The matrices  $E$  are **quantum effects**  $\iff \mathcal{D}_{\diamond,g}(1) \subseteq \mathcal{D}_{2E-I}(1)$
- The matrices  $E$  are **compatible quantum effects**  $\iff \mathcal{D}_{\diamond,g} \subseteq \mathcal{D}_{2E-I}$

At the intermediate levels  $1 \leq n \leq d$ ,  $\mathcal{D}_{\diamond,g}(n) \subseteq \mathcal{D}_{2E-I}(n)$  iff for all isometries  $V : \mathbb{C}^n \rightarrow \mathbb{C}^d$ , the compressed effects  $V^* E_i V$  are compatible.

Moreover, the compatibility region is equal to the set of inclusion constants of the matrix diamond:  $\forall g, d, \Gamma(g, d) = \Delta(g, d)$ .

The same results hold in the general (non-dichotomic) setting. One has to replace the matrix diamond by the **matrix jewel**  $\mathcal{D}_{\diamond,k}$  [BN20]

# Consequences

Many things are known about the matrix diamond

- For all  $g, d$ ,  $\frac{1}{2d}(1, 1, \dots, 1) \in \Delta(g, d)$  [HKMS19]
- For all  $g, d$ ,  $\text{QC}_g := \{s \in [0, 1]^g : \sum_i s_i^2 \leq 1\} \subseteq \Delta(g, d)$  [PSS18]

Many things are known about (in-)compatibility

- Many small  $g, d$  cases completely solved
- Approximate **quantum cloning** [Kay16]  $\implies$  compatibility

$$\text{Clone}(g, d) := \{s \in [0, 1]^g : \exists \text{ quantum channel } \Phi : \mathcal{M}_d \rightarrow \mathcal{M}_d^{\otimes g} \text{ s.t.} \\ \forall i \in [g], \quad \Phi_i(X) = s_i X + (1 - s_i) \frac{\text{Tr } X}{d}\}$$

- Connection with **tensor norms** [BN22b]

## Theorem ([BN18])

For all  $g$  and  $d \geq 2^{\lceil (g-1)/2 \rceil}$ ,  $\Gamma(g, d) = \Delta(g, d) = \text{QC}_g$

## Proof ideas

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# Inclusion of spectrahedra and (completely) positive maps

## Theorem ([HKM13])

Let  $A \in (\mathcal{M}_D^{sa}(\mathbb{C}))^g$ ,  $B \in (\mathcal{M}_d^{sa}(\mathbb{C}))^g$  such that  $\mathcal{D}_A(1)$  is bounded.

Then,  $\mathcal{D}_A(n) \subseteq \mathcal{D}_B(n)$  iff the unital linear map

$$\begin{aligned}\Phi : \text{span}\{I, A_1, \dots, A_g\} &\rightarrow \mathcal{M}_d^{sa}(\mathbb{C}) \\ A_i &\mapsto B_i\end{aligned}$$

is *n-positive*.

Sketch of the proof of the main theorem:

- Level 1: the extremal points of  $\mathcal{D}_{\diamond, g}(1)$  are  $\pm e_i$
- The inclusion  $\mathcal{D}_{\diamond, g} \subseteq \mathcal{D}_{2E-I}$  holds iff the unital map  $\Phi : I_2 \otimes \dots \otimes I_2 \otimes \text{diag}(1, -1) \otimes I_2 \otimes \dots \otimes I_2 \mapsto 2E_i - I_d$  is CP
- Arveson's extension theorem [Pau02, Theorem 6.2]:  $\Phi$  has a (completely) positive extension  $\tilde{\Phi}$  to  $\mathbb{R}^{2^g}$
- $C_\varepsilon := \tilde{\Phi}(\varepsilon)$  is a joint POVM for the  $E_i$ 's, where  $\{\varepsilon\}$  is a basis of  $\mathbb{R}^{2^g}$

# Maximally incompatible quantum effects

## Lemma ([Hru16])

For  $d = 2^k$ , there exist  $2k + 1$  *anti-commuting, self-adjoint, unitary* matrices  $F_1, \dots, F_{2k+1} \in \mathcal{U}_d$ . Moreover,  $2^k$  is the smallest dimension where such a  $(2k + 1)$ -tuple exists.

- For  $k = 0$ , take  $F_1^{(0)} := [1]$
- For  $k \geq 1$ , define  $F_i^{(k+1)} = \sigma_X \otimes F_i^{(k)} \ \forall i \in [2k + 1]$  and  $F_{2k+2}^{(k+1)} = \sigma_Y \otimes I_{2^k}$ ,  $F_{2k+3}^{(k+1)} = \sigma_Z \otimes I_{2^k}$
- These matrices satisfy, for all  $x \in \mathbb{R}_+^g$ ,  $\|\sum_{i=1}^g x_i F_i\|_\infty = \|x\|_2$ , and  $\|\sum_{i=1}^g x_i \bar{F}_i \otimes F_i\|_\infty = \|x\|_1$
- For  $d$  large enough, the maximally incompatible  $g$ -tuple of quantum effects in  $\mathcal{M}_d$  is given by  $E_i = (F_i + I_d)/2$

# The take-home slide

Measurement compatibility in QM  $\iff$  matrix diamond inclusion

- A measurement in QM (POVM):  $A_i \in \mathcal{M}_d(\mathbb{C})$ ,  $A_i \geq 0$ ,  $\sum_i A_i = I_d$
- Measurements  $(A_i), (B_j)$  are **compatible** if there exists a joint measurement  $(C_{ij})$  having marginals  $A, B$ :  $A_i = \sum_j C_{ij}$  and  $B_j = \sum_i C_{ij}$
- Free spectrahedra:  $\mathcal{D}_A := \bigsqcup_{n=1}^{\infty} \{X \in (\mathcal{M}_n^{sa})^g : \sum_{i=1}^g X_i \otimes A_i \leq I_{nd}\}$
- **Matrix diamond**:  $\mathcal{D}_{\diamond, g} = \bigsqcup_{n=1}^{\infty} \{X \in \mathcal{M}_n^g : \sum_{i=1}^g \varepsilon_i X_i \leq I_n, \forall \varepsilon \in \{\pm 1\}^g\}$

## Theorem

Let  $E \in (\mathcal{M}_d^{sa})^g$  be  $g$ -tuple of selfadjoint matrices. Then:

$$E \text{ are } \textbf{quantum effects} \iff \mathcal{D}_{\diamond, g}(1) \subseteq \mathcal{D}_{2E-I}(1)$$

$$E \text{ are } \textbf{compatible quantum effects} \iff \mathcal{D}_{\diamond, g} \subseteq \mathcal{D}_{2E-I}$$

For all  $g, d$ , compatibility region  $\Gamma(g, d) = \Delta(g, d)$  inclusion constants.

Steering inequality violation  $\iff$  matrix cube inclusion

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