

Some applications of free spectrahedra to quantum information theory

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33rd IWOTA, Krakow — September 10th 2022



What this talk is about

We establish a connection between two topics:

- **Incompatibility** of quantum measurements
- Inclusion of **free spectrahedra**, geometric objects which generalize polyhedra

The relation we put forward allows us to find the **most incompatible** quantum measurements [DFK19] in different settings. These objects are of great importance for the **foundations of quantum mechanics** due to their non-classical behavior.

We can also import quantum information techniques (such as **quantum cloning**) to the inclusion problem of free spectrahedra. The latter can be seen as one possible **positive semidefinite relaxation** of combinatorial optimization problems [HKMS19].

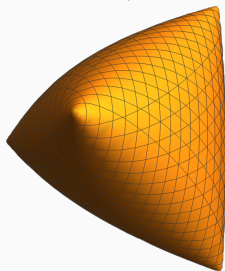
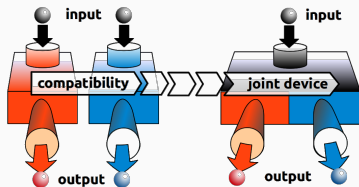
Talk outline

Incompatibility in QM

Free spectrahedra

Main results

Proof ideas



Incompatibility in QM

Quantum measurements. Compatibility

- Quantum states \rightsquigarrow **unit vectors in a Hilbert space**: $\psi \in \mathbb{C}^d$, $\|\psi\| = 1$
- Quantum measurements are modeled by **POVMs** (Positive Operator Valued Measures)

$$A_1, \dots, A_k \in M_d(\mathbb{C})^{\text{sa}}, \quad A_i \geq 0, \quad \sum_{i=1}^k A_i = I_d$$

- **Born's rule**: $\mathbb{P}[\text{outcome } i \text{ is observed}] = \langle \psi, A_i \psi \rangle$

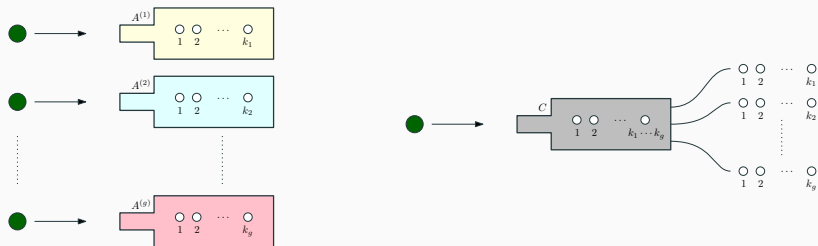
Definition

Two POVMs, $A = (A_1, \dots, A_k)$ and $B = (B_1, \dots, B_l)$, are called **compatible** if there exists a third POVM $C = (C_{ij})_{i \in [k], j \in [l]}$ such that

$$\forall i \in [k], \quad A_i = \sum_{j=1}^l C_{ij} \quad \text{and} \quad \forall j \in [l], \quad B_j = \sum_{i=1}^k C_{ij}$$

The definition generalizes to g -tuples of POVMs $A^{(1)}, \dots, A^{(g)}$, having respectively k_1, \dots, k_g outcomes, where the **joint** POVM C has outcome set $[k_1] \times \dots \times [k_g]$.

What does it mean?



- Compatible measurements can be simulated by a single joint measurement, by **classically post-processing** its outputs
- Compatibility appears also in the setting of quantum channels [HMZ16]

The compatibility game

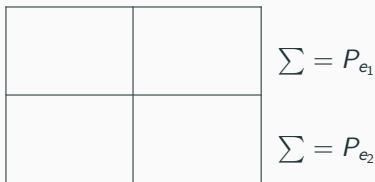
			$\Sigma = A_1$
			$\Sigma = A_2$
			$\Sigma = A_3$

$\Sigma = B_1$ $\Sigma = B_2$ $\Sigma = B_3$

We are given the Alice's $A = (A_1, A_2, A_3)$ and Bob's $B = (B_1, B_2, B_3)$ POVMs. **GOAL:** Fill in the 3×3 table with **positive semidefinite operators** (Marek ❤️) such that the row sums are the A 's and the column sums are the B 's

Two examples

- **Commuting** POVMs $[A_i, B_j] = 0$ are **compatible**: define $C_{ij} = A_i B_j = A_i^{1/2} B_j A_i^{1/2}$. Since the operators commute, $C_{ij} \geq 0$
Partial converse: if the POVM A is **projective** ($A_i^2 = A_i$), then A and B are compatible iff they commute
- Consider e the canonical basis of \mathbb{C}^2 and the “unbiased” basis $f = (f_1, f_2)$ with $f_1 = \begin{pmatrix} 1 \\ 1 \end{pmatrix} / \sqrt{2}$ and $f_2 = \begin{pmatrix} 1 \\ -1 \end{pmatrix} / \sqrt{2}$. The dichotomic (2-outcome) POVMs $A = (P_{e_1}, P_{e_2})$ and $B = (P_{f_1}, P_{f_2})$ are **incompatible**



$$\Sigma = P_{f_1} \quad \Sigma = P_{f_2}$$

Noisy POVMs

- POVMs can be made compatible by adding **noise**
- **Example:** dichotomic POVMs and white noise, $s \in [0, 1]$

$$(A_1, A_2) \mapsto s(A_1, A_2) + (1-s)\left(\frac{I}{2}, \frac{I}{2}\right) \quad \text{or} \quad A_i \mapsto sA_i + (1-s)\frac{I}{2}$$

- Taking $s = 1/2$ suffices to render any pair of dichotomic POVMs compatible [DFK19]
 - \rightsquigarrow define $C_{ij} := (A_i + B_j)/4$
 - \rightsquigarrow we have $\sum_j C_{ij} = (\sum_j A_i + \sum_j B_j)/4 = \frac{1}{2}A_i + \frac{1}{2}I$
- From now on, we focus on dichotomic (YES/NO) POVMs $(E, I - E)$ for operators $0 \leq E \leq I$ called **quantum effects**

Definition

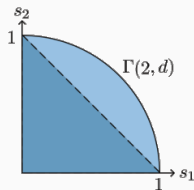
The **compatibility region** for g measurements on \mathbb{C}^d is the set

$$\Gamma(g, d) := \{s \in [0, 1]^g : \text{for all quantum effects } E_1, \dots, E_g \in \mathcal{M}_d(\mathbb{C}), \\ \text{the noisy versions } s_i E_i + (1 - s_i)I_d/2 \text{ are compatible}\}$$

Compatibility region

$\Gamma(g, d) := \{s \in [0, 1]^g : \text{for all quantum effects } E_1, \dots, E_g \in \mathcal{M}_d(\mathbb{C}),$
the noisy versions $s_i E_i + (1 - s_i)I_d/2$ are **compatible**\}

- The set $\Gamma(g, d)$ is convex
- For all $i \in [g]$, $e_i \in \Gamma(g, d)$: every measurement is compatible with $g - 1$ trivial measurements
- For $d \geq 2$, $(1, 1, \dots, 1) \notin \Gamma(g, d)$: there exist incompatible measurements
- For all $d \geq 2$, $\Gamma(2, d)$ is a quarter-circle [BH08, BHSS13]
- Generally speaking, the set $\Gamma(g, d)$ tells us how **robust** (to noise) is the incompatibility of g dichotomic measurements on \mathbb{C}^d



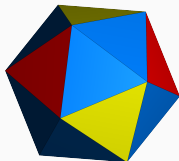
GOAL: compute the set $\Gamma(g, d)$ for all g, d

Free spectrahedra

Free spectrahedra

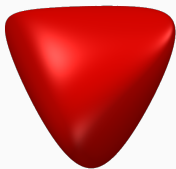
- A **polyhedron** is defined as the intersection of half-spaces

$$\{x \in \mathbb{R}^g : \langle h_i, x \rangle \leq 1, \quad \forall i \in [k]\}$$



- A **spectrahedron** is given by PSD constraints:
for $A = (A_1, \dots, A_g) \in (\mathcal{M}_d^{sa})^g$

$$\mathcal{D}_A(1) := \{x \in \mathbb{R}^g : \sum_{i=1}^g x_i A_i \leq I_d\}$$



- $\mathcal{D}_{(\sigma_X, \sigma_Y, \sigma_Z)}(1) = \{(x, y, z) \in \mathbb{R}^3 : x\sigma_X + y\sigma_Y + z\sigma_Z \leq I_2\} = \text{Bloch ball}$
- A **free spectrahedron** is the matricization of a spectrahedron [Vin14]

$$\mathcal{D}_A := \bigsqcup_{n=1}^{\infty} \mathcal{D}_A(n) \quad \text{with} \quad \mathcal{D}_A(n) := \{X \in (\mathcal{M}_n^{sa})^g : \sum_{i=1}^g X_i \otimes A_i \leq I_{nd}\}$$

Examples: the cube and the diamond

The **matrix cube** is the free spectrahedron defined by

$$\mathcal{D}_{\square, g} := \bigsqcup_{n=1}^{\infty} \{X \in (\mathcal{M}_n^{sa})^g : \|X_i\| \leq 1, \quad \forall i \in [g]\}$$

- At level one, $\mathcal{D}_{\square, g}(1)$ is the unit ball of the ℓ^∞ norm on \mathbb{R}^g
- As a free spectrahedron, it is defined by $2g \times 2g$ diagonal matrices $\mathcal{D}_{\square, g} = \mathcal{D}_{K_1, \dots, K_g}$, with $K_i = \text{diag}(e_i) \oplus \text{diag}(-e_i)$

The **matrix diamond** is the free spectrahedron defined by

$$\mathcal{D}_{\diamond, g} := \bigsqcup_{n=1}^{\infty} \{X \in (\mathcal{M}_n^{sa})^g : \sum_{i=1}^g \varepsilon_i X_i \leq I_n, \quad \forall \varepsilon \in \{\pm 1\}^g\}$$

- At level one, $\mathcal{D}_{\diamond, g}(1)$ is the unit ball of the ℓ^1 norm on \mathbb{R}^g
- As a free spectrahedron, it is defined by $2^g \times 2^g$ diagonal matrices $\mathcal{D}_{\diamond, g} = \mathcal{D}_{L_1, \dots, L_g}$, with $L_i = I_2 \otimes \cdots \otimes I_2 \otimes \text{diag}(1, -1) \otimes I_2 \otimes \cdots \otimes I_2$

Spectrahedral inclusion

- Consider two free spectrahedra defined by g -tuples of matrices (A_1, \dots, A_g) and (B_1, \dots, B_g)
- We say that \mathcal{D}_A is **contained** in \mathcal{D}_B and we write $\mathcal{D}_A \subseteq \mathcal{D}_B$ if, for all $n \geq 1$, $\mathcal{D}_A(n) \subseteq \mathcal{D}_B(n)$
- Clearly, $\mathcal{D}_A \subseteq \mathcal{D}_B \implies \mathcal{D}_A(1) \subseteq \mathcal{D}_B(1)$. For the converse implication to hold, one may need to shrink \mathcal{D}_A ...

Definition

For a free spectrahedron \mathcal{D}_A , we define its set of **inclusion constants** as

$$\Delta_A(g, d) := \{s \in [0, 1]^g : \text{for all } g\text{-tuples } B_1, \dots, B_g \in \mathcal{M}_d(\mathbb{C})^{\text{sa}}, \\ \mathcal{D}_A(1) \subseteq \mathcal{D}_B(1) \implies s \cdot \mathcal{D}_A \subseteq \mathcal{D}_B\}$$

- The inclusion constants for the matrix cube play an important role in combinatorial optimization [HKMS19]
- We shall be concerned with the inclusion set for the **matrix diamond**, which we denote by $\Delta(g, d)$; it is the same as the one of the matrix cube [BN22a]

Main results

Compatibility in QM \iff matrix diamond inclusion

To a g -tuple of selfadjoint matrices $E \in (\mathcal{M}_d^{sa})^g$, we associate the free spectrahedron defined by the matrices $2E_i - I_d$:

$$\mathcal{D}_{2E-I} := \bigsqcup_{n=1}^{\infty} \{X \in (\mathcal{M}_n^{sa})^g : \sum_{i=1}^g X_i \otimes (2E_i - I_d) \leq I_{nd}\}$$

Theorem ([BN18])

Let $E \in (\mathcal{M}_d^{sa})^g$ be g -tuple of selfadjoint matrices. Then:

- The matrices E are **quantum effects** $\iff \mathcal{D}_{\diamond,g}(1) \subseteq \mathcal{D}_{2E-I}(1)$
- The matrices E are **compatible quantum effects** $\iff \mathcal{D}_{\diamond,g} \subseteq \mathcal{D}_{2E-I}$

At the intermediate levels $1 \leq n \leq d$, $\mathcal{D}_{\diamond,g}(n) \subseteq \mathcal{D}_{2E-I}(n)$ iff for all isometries $V : \mathbb{C}^n \rightarrow \mathbb{C}^d$, the compressed effects $V^* E_i V$ are compatible.

Moreover, the compatibility region is equal to the set of inclusion constants of the matrix diamond: $\forall g, d, \Gamma(g, d) = \Delta(g, d)$.

The same results hold in the general (non-dichotomic) setting. One has to replace the matrix diamond by the **matrix jewel** $\mathcal{D}_{\diamond,k}$ [BN20]

Consequences

Many things are known about the matrix diamond

- For all g, d , $\frac{1}{2d}(1, 1, \dots, 1) \in \Delta(g, d)$ [HKMS19]
- For all g, d , $\text{QC}_g := \{s \in [0, 1]^g : \sum_i s_i^2 \leq 1\} \subseteq \Delta(g, d)$ [PSS18]

Many things are known about (in-)compatibility

- Many small g, d cases completely solved
- Approximate **quantum cloning** [Kay16] \implies compatibility

$$\text{Clone}(g, d) := \{s \in [0, 1]^g : \exists \text{ quantum channel } \Phi : \mathcal{M}_d \rightarrow \mathcal{M}_d^{\otimes g} \text{ s.t.} \\ \forall i \in [g], \quad \Phi_i(X) = s_i X + (1 - s_i) \frac{\text{Tr } X}{d}\}$$

- Connection with **tensor norms** [BN22b]

Theorem ([BN18])

For all g and $d \geq 2^{\lceil (g-1)/2 \rceil}$, $\Gamma(g, d) = \Delta(g, d) = \text{QC}_g$

Proof ideas

Inclusion of spectrahedra and (completely) positive maps

Theorem ([HKM13])

Let $A \in (\mathcal{M}_D^{sa}(\mathbb{C}))^g$, $B \in (\mathcal{M}_d^{sa}(\mathbb{C}))^g$ such that $\mathcal{D}_A(1)$ is bounded. Then, $\mathcal{D}_A(n) \subseteq \mathcal{D}_B(n)$ iff the unital linear map

$$\begin{aligned}\Phi : \text{span}\{I, A_1, \dots, A_g\} &\rightarrow \mathcal{M}_d^{sa}(\mathbb{C}) \\ A_i &\mapsto B_i\end{aligned}$$

is *n-positive*.

Sketch of the proof of the main theorem:

- Level 1: the extremal points of $\mathcal{D}_{\diamond, g}(1)$ are $\pm e_i$
- The inclusion $\mathcal{D}_{\diamond, g} \subseteq \mathcal{D}_{2E-I}$ holds iff the unital map $\Phi : I_2 \otimes \dots \otimes I_2 \otimes \text{diag}(1, -1) \otimes I_2 \otimes \dots \otimes I_2 \mapsto 2E_i - I_d$ is CP
- Arveson's extension theorem [Pau02, Theorem 6.2]: Φ has a (completely) positive extension $\tilde{\Phi}$ to \mathbb{R}^{2^g}
- $C_\varepsilon := \tilde{\Phi}(\varepsilon)$ is a joint POVM for the E_i 's, where $\{\varepsilon\}$ is a basis of \mathbb{R}^{2^g}

Maximally incompatible quantum effects

Lemma ([Hru16])

For $d = 2^k$, there exist $2k + 1$ *anti-commuting, self-adjoint, unitary* matrices $F_1, \dots, F_{2k+1} \in \mathcal{U}_d$. Moreover, 2^k is the smallest dimension where such a $(2k + 1)$ -tuple exists.

- For $k = 0$, take $F_1^{(0)} := [1]$
- For $k \geq 1$, define $F_i^{(k+1)} = \sigma_X \otimes F_i^{(k)} \forall i \in [2k + 1]$ and $F_{2k+2}^{(k+1)} = \sigma_Y \otimes I_{2^k}$, $F_{2k+3}^{(k+1)} = \sigma_Z \otimes I_{2^k}$
- These matrices satisfy, for all $x \in \mathbb{R}_+^g$, $\|\sum_{i=1}^g x_i F_i\|_\infty = \|x\|_2$, and $\|\sum_{i=1}^g x_i \bar{F}_i \otimes F_i\|_\infty = \|x\|_1$
- For d large enough, the maximally incompatible g -tuple of quantum effects in \mathcal{M}_d is given by $E_i = (F_i + I_d)/2$

The take-home slide

Measurement compatibility in QM \iff matrix diamond inclusion

- A measurement in QM (POVM): $A_i \in \mathcal{M}_d(\mathbb{C})$, $A_i \geq 0$, $\sum_i A_i = I_d$
- Measurements $(A_i), (B_j)$ are **compatible** if there exists a joint measurement (C_{ij}) having marginals A, B : $A_i = \sum_j C_{ij}$ and $B_j = \sum_i C_{ij}$
- Free spectrahedra: $\mathcal{D}_A := \bigsqcup_{n=1}^{\infty} \{X \in (\mathcal{M}_n^{sa})^g : \sum_{i=1}^g X_i \otimes A_i \leq I_{nd}\}$
- **Matrix diamond**: $\mathcal{D}_{\diamond, g} = \bigsqcup_{n=1}^{\infty} \{X \in \mathcal{M}_n^g : \sum_{i=1}^g \varepsilon_i X_i \leq I_n, \forall \varepsilon \in \{\pm 1\}^g\}$

Theorem

Let $E \in (\mathcal{M}_d^{sa})^g$ be g -tuple of selfadjoint matrices. Then:

$$E \text{ are } \textit{quantum effects} \iff \mathcal{D}_{\diamond, g}(1) \subseteq \mathcal{D}_{2E-I}(1)$$

$$E \text{ are } \textit{compatible quantum effects} \iff \mathcal{D}_{\diamond, g} \subseteq \mathcal{D}_{2E-I}$$

For all g, d , compatibility region $\Gamma(g, d) = \Delta(g, d)$ inclusion constants.

Steering inequality violation \iff matrix cube inclusion

References

- [BH08] Paul Busch and Teiko Heinosaari.
Approximate joint measurements of qubit observables.
Quantum Information & Computation, 8(8):797–818, 2008.
- [BHSS13] Paul Busch, Teiko Heinosaari, Jussi Schultz, and Neil Stevens.
Comparing the degrees of incompatibility inherent in probabilistic physical theories.
EPL (Europhysics Letters), 103(1):10002, 2013.
- [BN18] Andreas Bluhm and Ion Nechita.
Joint measurability of quantum effects and the matrix diamond.
Journal of Mathematical Physics, 59(11):112202, 2018.
- [BN20] Andreas Bluhm and Ion Nechita.
Compatibility of quantum measurements and inclusion constants for the matrix jewel.
SIAM Journal on Applied Algebra and Geometry, 4(2):255–296, 2020.
- [BN22a] Andreas Bluhm and Ion Nechita.
Maximal violation of steering inequalities and the matrix cube.
Quantum, 6:656, 2022.
- [BN22b] Andreas Bluhm and Ion Nechita.
A tensor norm approach to quantum compatibility.
Journal of Mathematical Physics, 63(6):062201, 2022.
- [DFK19] Sébastien Designolle, Máté Farkas, and Jędrzej Kaniewski.
Incompatibility robustness of quantum measurements: a unified framework.
New Journal of Physics, 21(11):113053, 2019.
- [HKM13] J. William Helton, Igor Klep, and Scott McCullough.
The matricial relaxation of a linear matrix inequality.
Mathematical Programming, 138(1-2):401–445, 2013.
- [HKMS19] J. William Helton, Igor Klep, Scott McCullough, and Markus Schweighofer.
Dilations, linear matrix inequalities, the matrix cube problem and beta distributions.
Memoirs of the American Mathematical Society, 257(1232), 2019.
- [HMZ16] Teiko Heinosaari, Takayuki Miyadera, and Mário Ziman.
An invitation to quantum incompatibility.
Journal of Physics A: Mathematical and Theoretical, 49(12):123001, 2016.
- [Hru16] Pavel Hrušeš.
On families of anticommuting matrices.
Linear Algebra and its Applications, 493:494–507, 2016.
- [Kay16] Alastair Kay.
Optimal universal quantum cloning: Asymmetries and fidelity measures.
Quantum Information and Computation, 16(11 & 12):0991–1028, 2016.
- [Pau02] Vern Paulsen.
Completely bounded maps and operator algebras, volume 78.
Cambridge University Press, 2002.
- [PSS18] Benjamin Passer, Orr Moshe Shalit, and Baruch Solel.
Minimal and maximal matrix convex sets.
Journal of Functional Analysis, 274:3197–3253, 2018.
- [Vin14] Cynthia Vinzant.
What is... a spectrahedron.
Notices Amer. Math. Soc., 61(5):492–494, 2014.